

The 40th JKPS:

Longitudinal-to-transverse mapping for femtosecond electron bunch length measurement

Dao Xiang and Yuantao Ding

SLAC National Accelerator Laboratory, Menlo Park, California 94025, USA

(Received 26 April 2010; published 1 September 2010)

A longitudinal-to-transverse mapping technique is proposed to measure the length and temporal profile of ultrashort electron bunches. In this scheme a special chicane and a radio-frequency deflecting cavity are used to transform the beam's longitudinal distribution into angular distribution which is further converted to transverse distribution after a parallel-to-point imaging beam line. With this technique, the temporal profile of the electron beam is exactly mapped to the transverse profile. This makes it possible to measure ultrashort electron bunch length with a resolution well beyond 1 femtosecond.

DOI: [10.1103/PhysRevSTAB.13.094001](https://doi.org/10.1103/PhysRevSTAB.13.094001)

PACS numbers: 29.27.Fh, 41.60.Cr

26. 01. 23

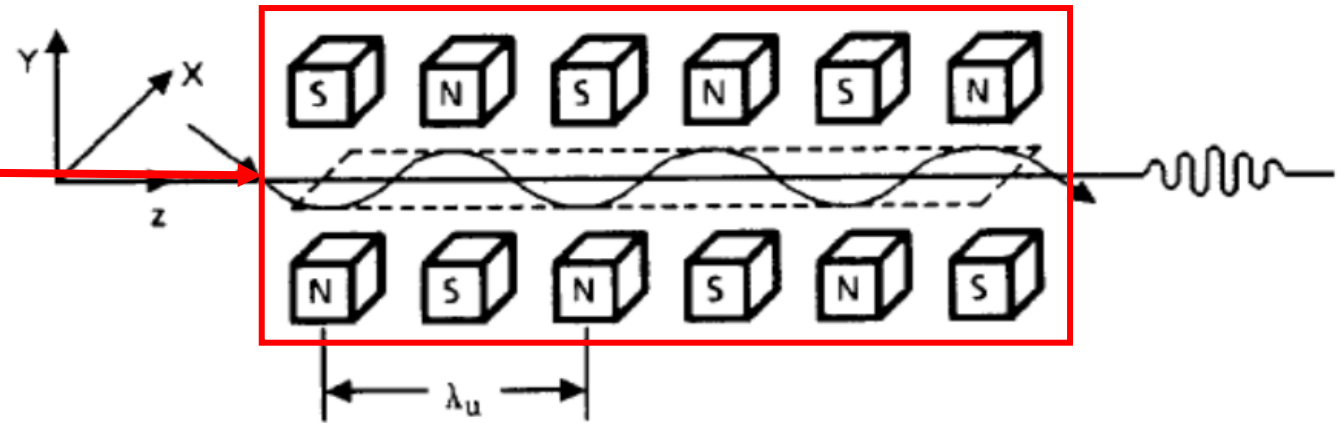
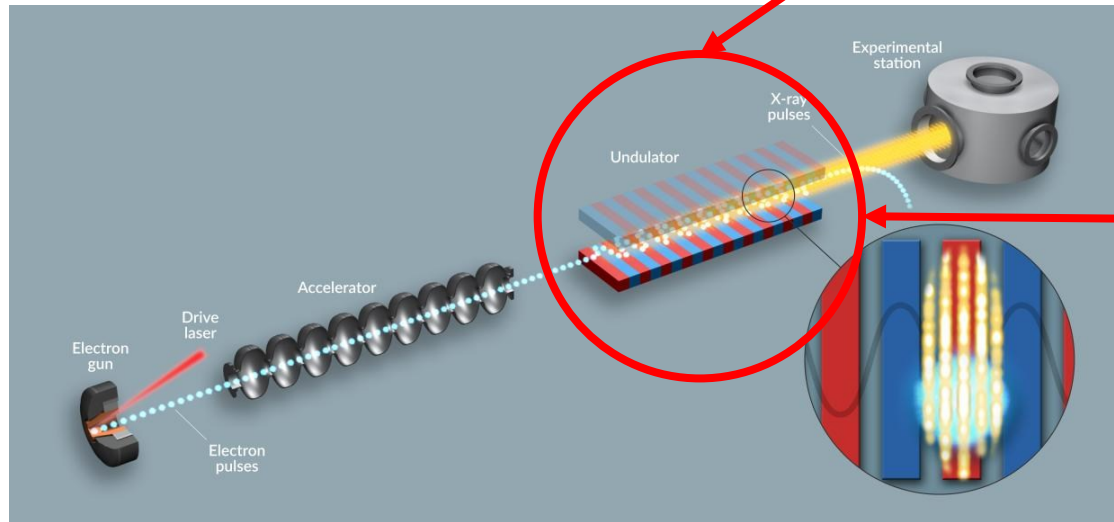
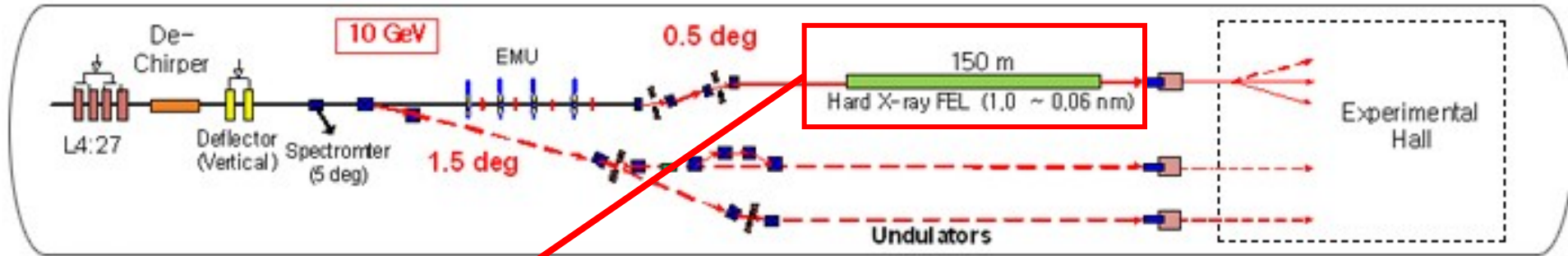
Geunwoo Kim

POSTECH



Background: Free Electron Laser (FEL)

- Schematic of FEL beamline



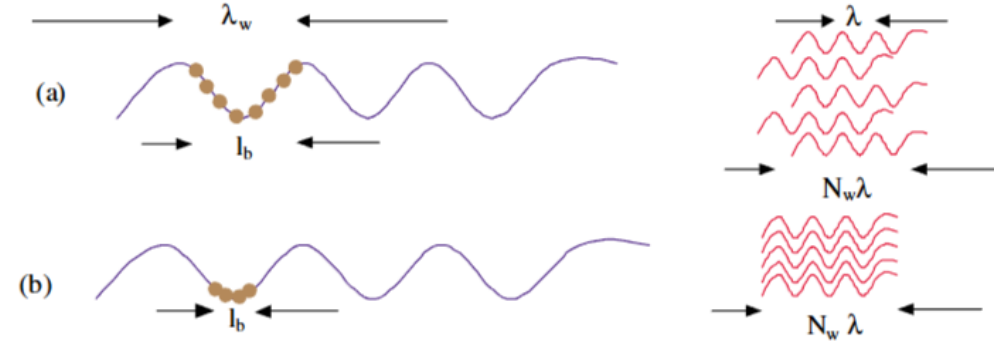
Background: Free Electron Laser (FEL)

- Requirements for FEL performance

1. Short pulse (=short length beam): High peak current → interests in this paper (How to measure it?)
2. Low energy spread is essential for stable FEL gain.

Gain length: $L_g = \frac{\lambda_u}{4\pi\sqrt{3}\cdot\rho}$

Pierce parameter: $\rho^3 = \frac{I}{4I_A} \frac{\gamma K_u^2 [JJ]^2 (K_u^2)}{\left(1 + \frac{K_u^2}{2}\right)^2 (k_r \sigma_x)^2}, \rightarrow \rho^3 \propto I / \sigma_x^2$

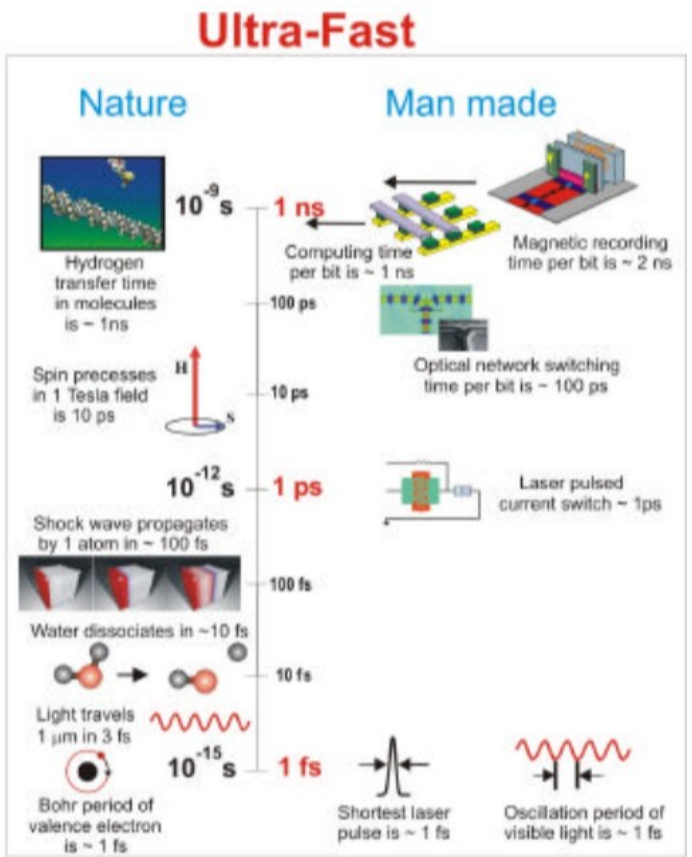
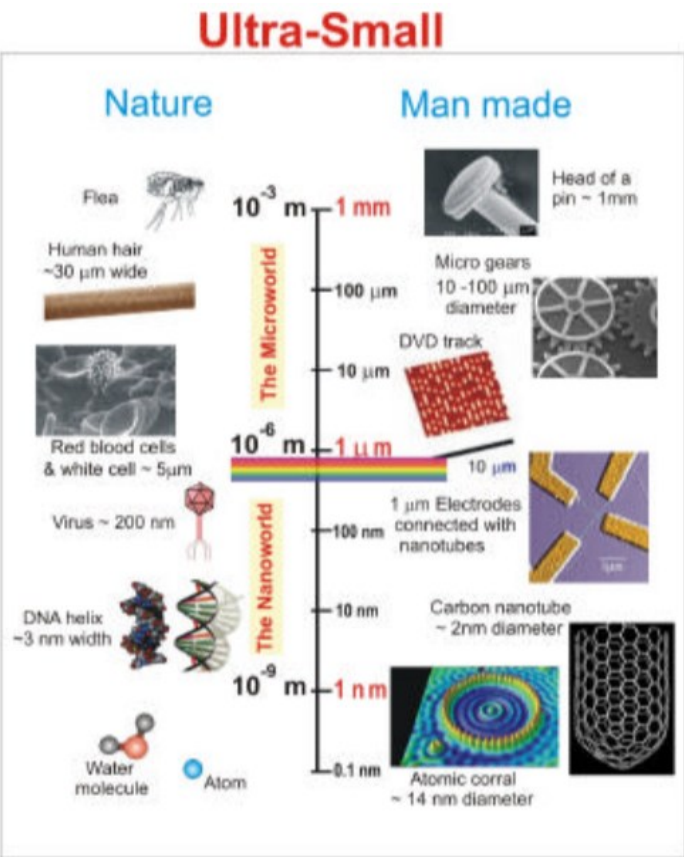
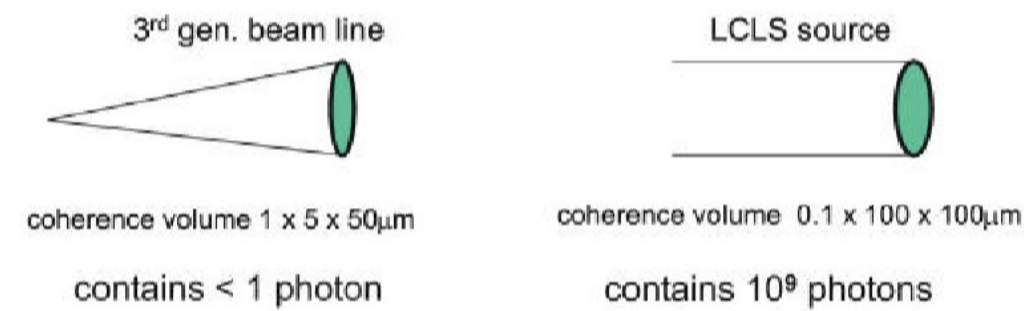


“Increasing the peak current shortens the gain length.”

Background: Free Electron Laser (FEL)

- Application

	3 rd Gen.	SASE-FREE-ELECTRON LASER	Short pulse SASE-FREE-ELECTRON LASER
Wavelength range, nm	1-0.1		
Emittance, nm rad	2	0.03	0.03
Pulse length, ps	15-30	0.06	0.01
Average brightness	10^{20}	10^{22}	10^{21}
Peak brightness	10^{23}	10^{33}	10^{33}
Peak power, W	10^3	10^{10}	10^{10}



Background: Dipole Magnet

Main feature

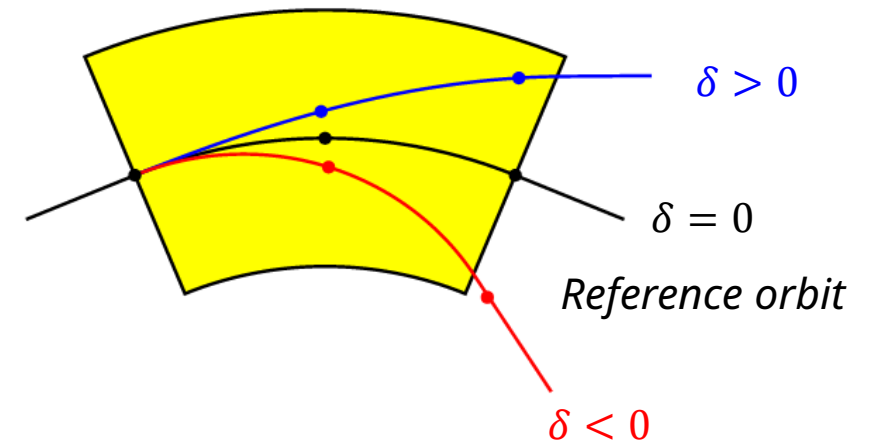
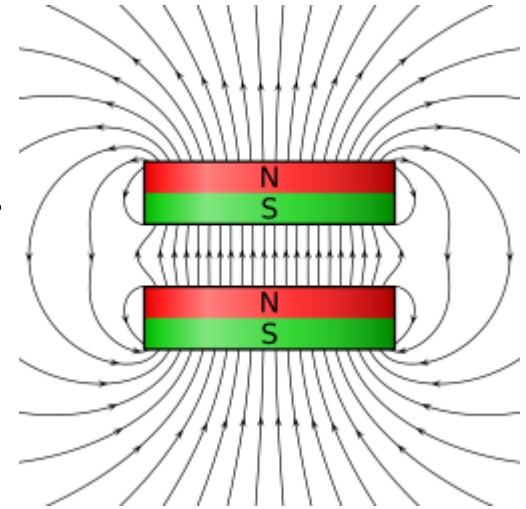
- Generates a uniform magnetic field to bend charged particles.
- Bending angle θ depends on momentum $\rightarrow$ energy-dependent dispersion.

$$\theta \approx 0.3 \cdot \frac{B [\text{T}] \cdot L [\text{m}]}{E + \Delta E [\text{GeV}]}$$

- Maps energy deviation (δ) into horizontal position.

$$\delta = \frac{p - p_0}{p_0}$$

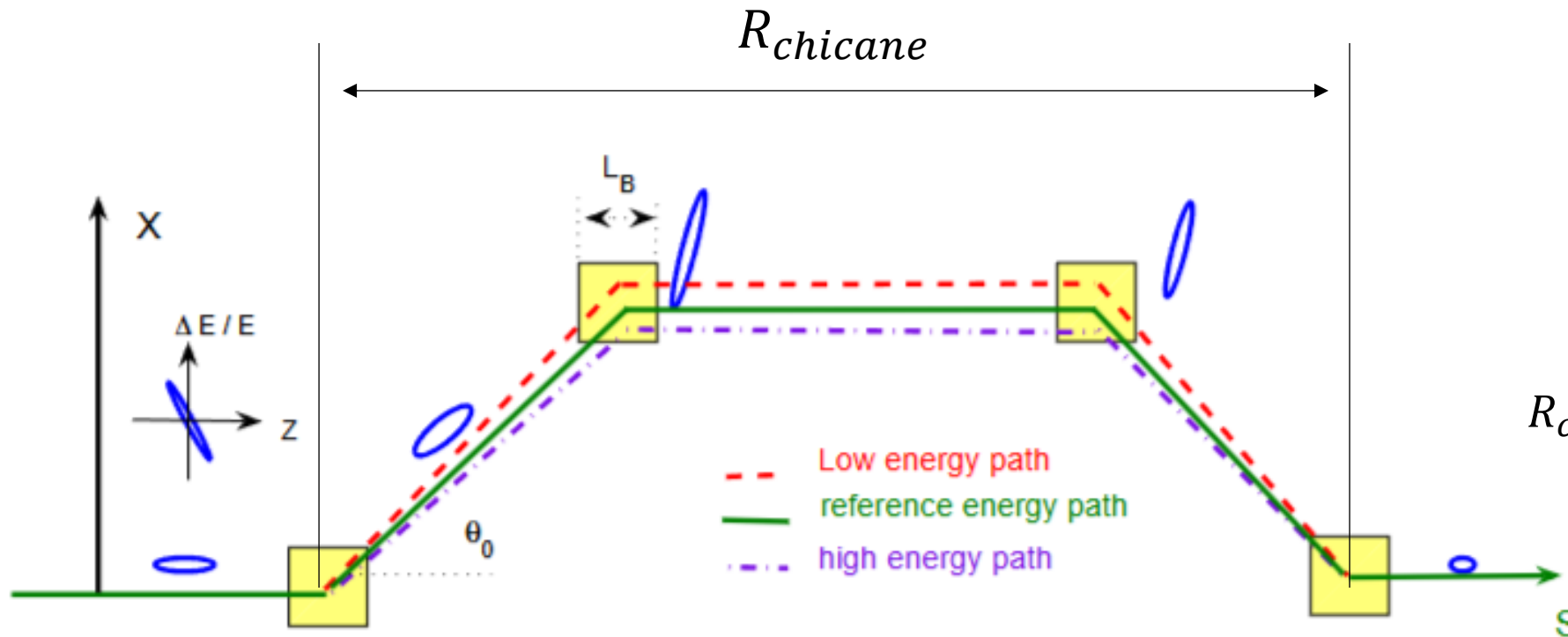
$$x(s) = x_{\beta(s)} + \mathbf{D}(s)\delta$$



Background: Chicane (usually bunch compressor)

Main feature

- Made of 4 dipole magnets.
- *Used to shorten the electron bunch length by exploiting $\delta - x$ coupling.*
- $z_f = z_i + \mathbf{R}_{34}\delta$ where $R_{56} = \int \frac{D(s)}{\rho(s)} ds$



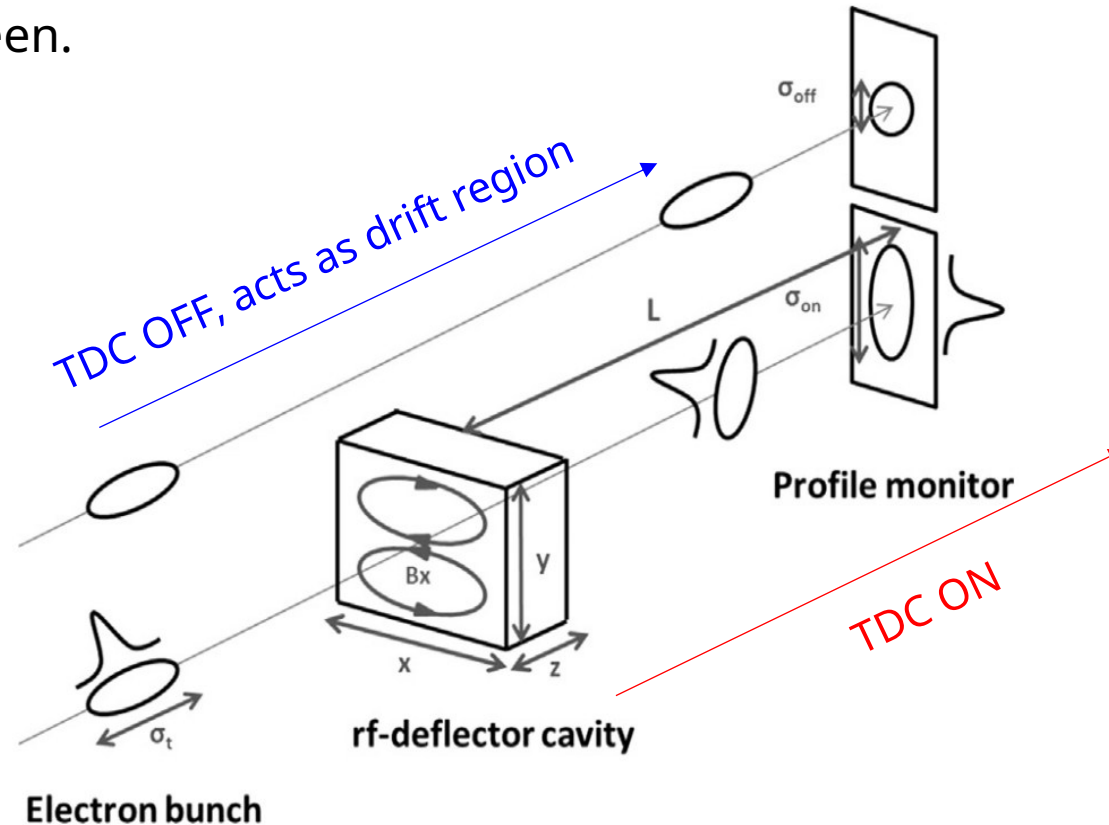
$$\vec{g}_0 = (x, x', z, \delta)^T$$

$$R_{chicane} = \begin{bmatrix} 1 & L & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & \mathbf{R}_{34} \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

Background: Transverse Deflecting Structure (TDS)

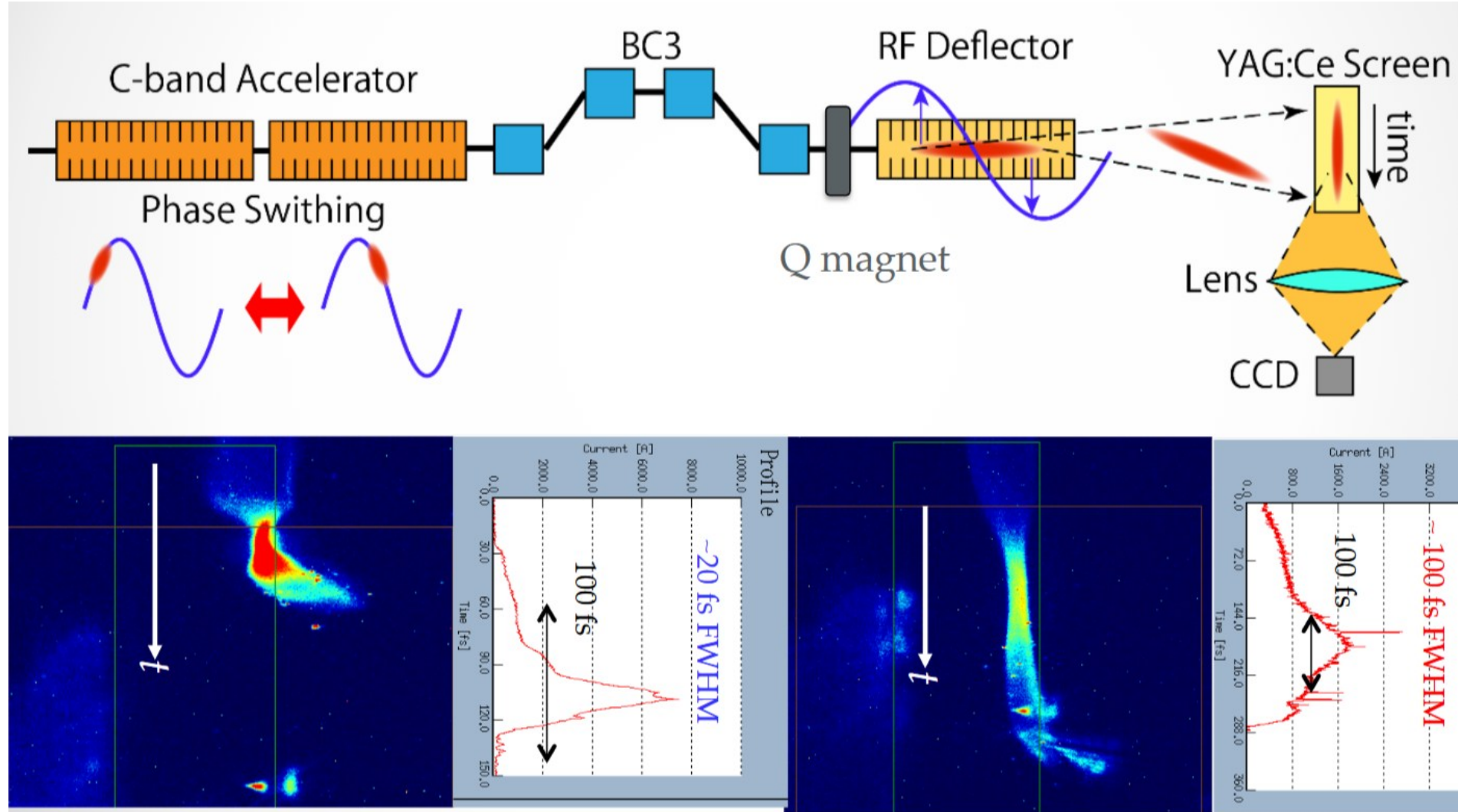
- **Basic principle of TDS**

- A type of the longitudinal beam profile monitor.
- Provides a transverse kick via RF fields ($\vec{F} = q (\vec{E} + \vec{v} \times \vec{B})$).
- Map $z \rightarrow x$ or y on screen.



Existing Method: Bunch Length Measurement

- RF deflector + YAG screen measure bunch length ($\sim 10\text{--}20$ fs).



Introduction: Reported Bunch Length at LCLS

Motivations

- LCLS achieves bunch lengths of ~ 3 fs.
- Achieved 10-20 fs resolution through TDS (2010).
- Such short pulses are beyond TDS diagnostics.
- Plus, all proposed method depends on $\sigma_{intrinsic}$.

→ How to measure sub-fs scale beam length?

Parameter ^a	20 pC	250 pC	1 nC	Unit
UV laser energy on cathode	1.5	20	250	μJ
UV spot diameter on cathode	0.6	1.2	2	mm
UV pulse duration (fwhm)	4.0	6.5	10	ps
Injector bunch length (rms)	1.3	2.5	2.8	ps
Initial peak current	5	30	100	A
Injector slice emittance	0.14	0.6	1.0	μm
Injector projected emittance	0.20	0.7	1.2	μm
Final bunch length (rms)	~ 3	~ 30	80	fs
Final peak current	~ 3	~ 3	3.4	kA
Final projected emittance	0.4	1.0	1.5	μm
FEL pulse duration (fwhm) ^b	~ 2	~ 60	230	fs
FEL peak power ^b	~ 400	~ 20	~ 10	GW

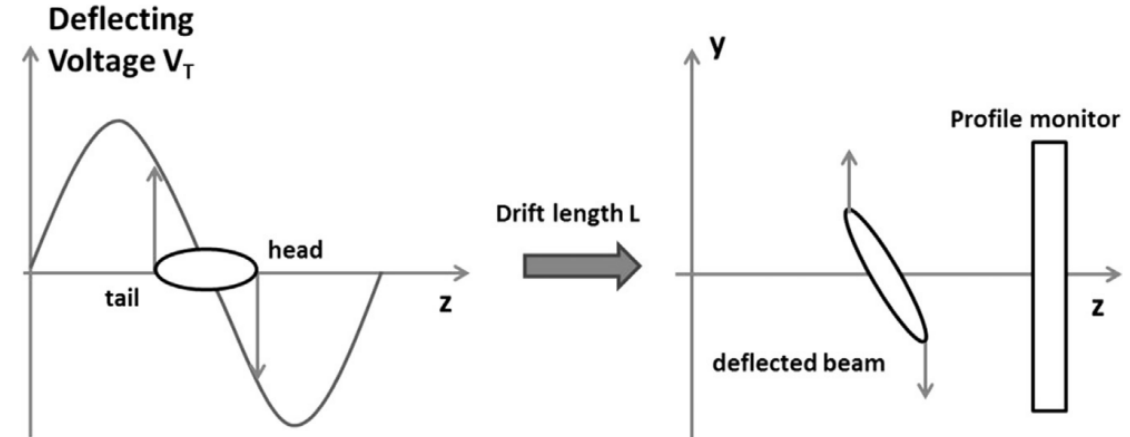
Resolution Limit of TDS

- Temporal resolution of TDS

- $$\sigma_{t,res} = \frac{p_z c}{e} \cdot \sigma_{x0} \cdot \frac{1}{V_T \omega_{RFL}}$$

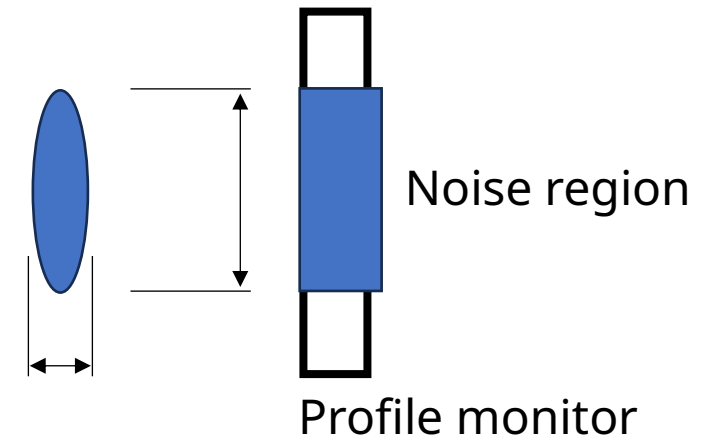
- $$R_D R_{TDS} = \begin{bmatrix} 1 & L & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & k & 0 \\ 0 & 0 & 1 & 0 \\ k & 0 & 0 & 1 \end{bmatrix} \text{ for } \beta \approx 1$$

- $$R_{final} = \begin{bmatrix} 1 & L & kL & 0 \\ 0 & 1 & k & 0 \\ 0 & 0 & 1 & 0 \\ k & 0 & 0 & 1 \end{bmatrix}$$



If $\sigma_x > \sigma_z \dots$

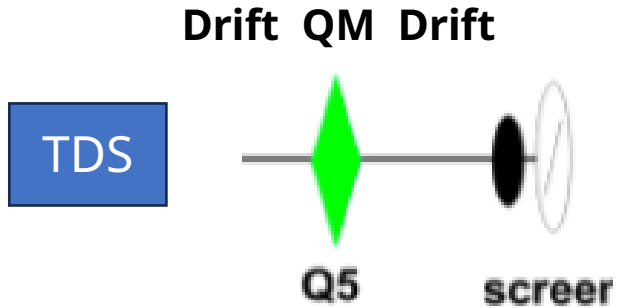
$$\begin{aligned} \rightarrow x_f &= \text{Noise } x_i + \text{Signal } Lx'_i + kLz_i \\ \rightarrow \sigma_x^2 &= \sigma_{x0}^2 + L^2 k^2 \sigma_{z0}^2 \end{aligned}$$



Method: elimination x_i

Parallel-to-Point optics

- D-Q-D setup enables removal of x_i .
- Setting the focal length to $f = L$ yields $R_{11} = 0$
- Transfer map:



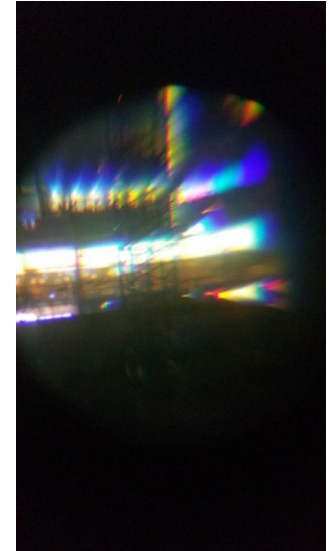
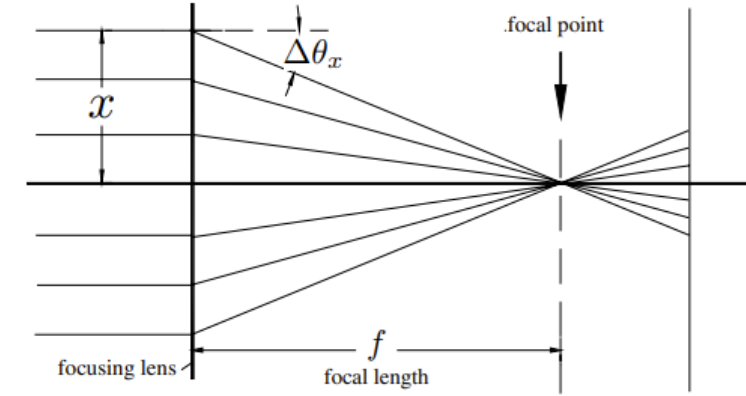
$$R_D R_{QF} R_D = \begin{bmatrix} 1 - \frac{L}{f} & L(2 - \frac{L}{f}) \\ -\frac{1}{f} & 1 - \frac{L}{f} \end{bmatrix}$$

$$R_{PtP} = \begin{bmatrix} 0 & f \\ -\frac{1}{f} & 0 \end{bmatrix} \text{ for } f = L.$$

$$R_{PtP} R_{TDS} = \begin{bmatrix} 0 & f & 0 & 0 \\ -\frac{1}{f} & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & k & 0 \\ 0 & 0 & 1 & 0 \\ k & 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} 0 & f & kf & 0 \\ 0 & 1 & k & 0 \\ 0 & 0 & 1 & 0 \\ k & 0 & 0 & 1 \end{bmatrix} \rightarrow \text{Removal of } R_{11}$$

$$x_f = \cancel{x_i} + Lx'_i + kLz_i$$

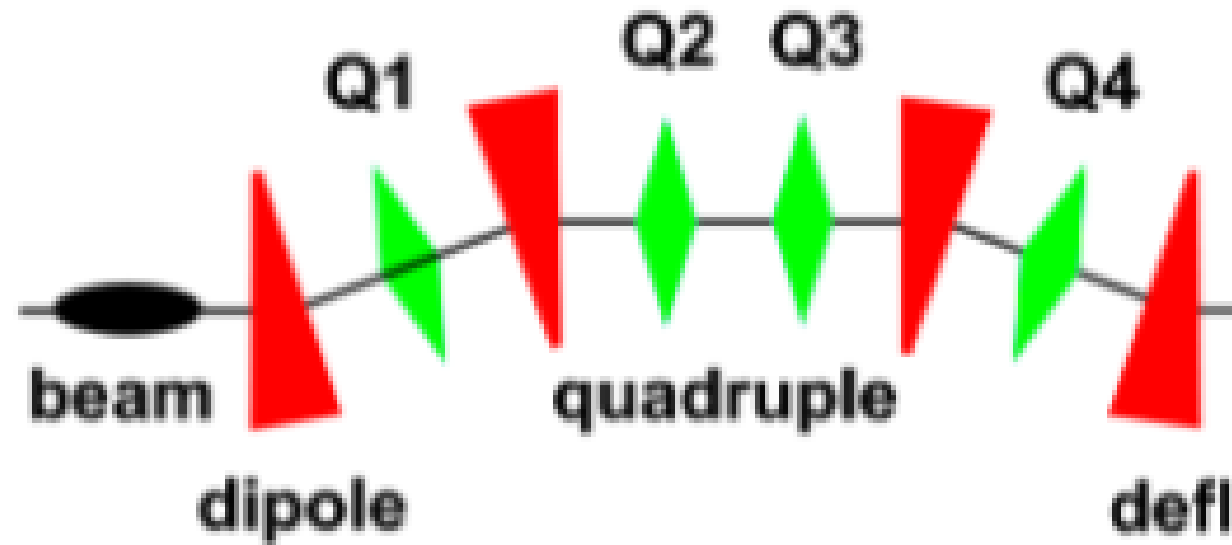
$$\vec{F}_{QM} = q\vec{v} \times \vec{B} = qv_z K \begin{bmatrix} -x \\ y \\ 0 \end{bmatrix}$$



Method: Special chicane to remove x'_i

Special chicane for eliminating x'_i

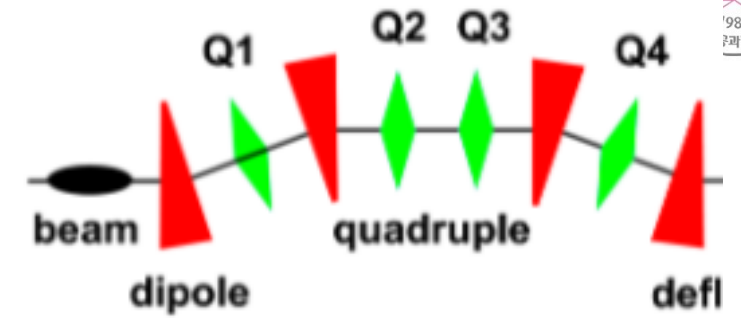
- Insert symmetric QM inside the chicane.
- *Control the dispersion* and match with the TDS strength k .
- Using $-I$ transformer (mirror-symmetric lattice with $\Delta\phi = \pi$)



Method: Special chicane to remove x'_i

Special chicane with $D \neq 0$, $D' = 0$

- I transformer sets $R_{11} = -1$, $R_{22} = -1$
- Constraint via QM
 - $R_{31} = 0$ (Independent between $x - z$)
 - $R_{33} = 1$ (No bunch compression)
 - $R_{34} = 0$ (Isochronous chicane)



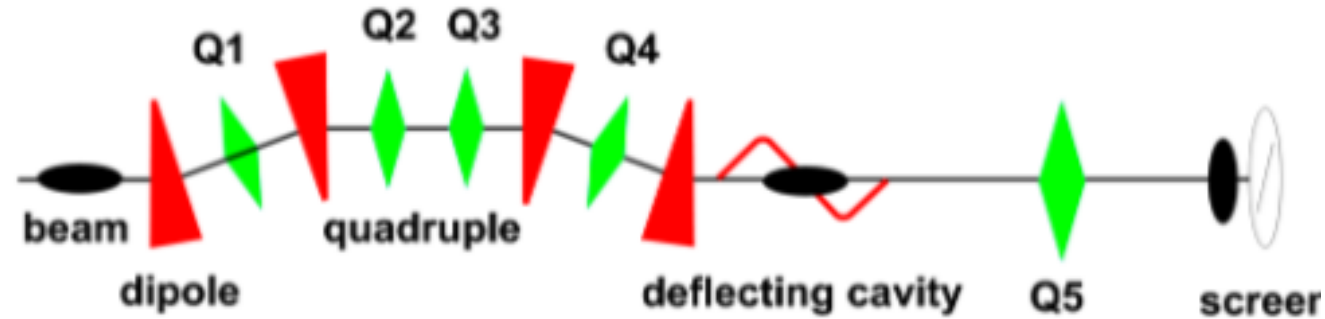
$$\begin{bmatrix} x_c \\ x'_c \\ z_c \\ \delta_c \end{bmatrix} = \begin{bmatrix} -1 & R_{12} & 0 & \eta \\ R_{21} & -1 & 0 & 0 \\ 0 & R_{32} & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} x_i \\ x'_i \\ z_i \\ \delta_i \end{bmatrix} \quad D \neq 0$$

- Symplectic condition $R^T S R = S$ where $S = \begin{bmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{bmatrix} \rightarrow (R_{12}R_{24} - R_{22}R_{14}) + (R_{32}R_{44} - R_{42}R_{34}) = 0$
 $(0 + \eta) + (R_{32} \cdot 1 - 0) = 0$
 $\therefore R_{32} = -\eta, R_{21} = 0$

Method: Special chicane to remove x'_i

Calculation whole transfer map

- If one sets to dispersion $k = 1/\eta$, possible to remove initial x'_i .



$$R_{TDS}R_c = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & k & 0 \\ 0 & 0 & 1 & 0 \\ k & 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} -1 & R_{12} & 0 & \eta \\ 0 & -1 & 0 & 0 \\ 0 & -\eta & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} -1 & R_{12} & 0 & \eta \\ 0 & -(1 + \eta k) & k & 0 \\ 0 & -\eta & 1 & 0 \\ -k & kR_{12} & 0 & 1 + \eta k \end{bmatrix}$$

$$R_{total} = R_{PtP}R_{TDS}R_c = \begin{bmatrix} 0 & 0 & fk & 0 \\ 1/f & -R_{12}/f & 0 & 1/fk \\ 0 & 1/k & 1 & 0 \\ -k & kR_{12} & 0 & 0 \end{bmatrix} \text{ for } k = 1/\eta.$$

$$x_f = \cancel{x_i} + L\cancel{x'_i} + kLz_i$$

Effective magnification ratio $M = kL$

Verification: ELEGANT simulation

ELEGANT simulation for 2nd -order effects

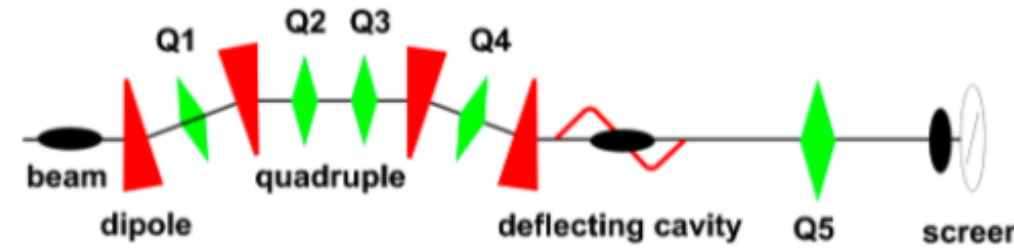
- Higher-order terms (T_{266}, T_{566}) must be considered.
- They leave a residual σ_{x0} .
 - Nonlinear effects cause additional spread and worse resolution.

$$x'_1 \approx kz_1 + T_{266}\delta_0^2$$

$$z_1 \approx z_0 + T_{566}\delta_0^2$$

$$\Delta t \approx \sqrt{\frac{2T_{266}^2}{k^2} + 2T_{566}^2} \cdot \frac{\delta_0^2}{c}$$

- Initial transverse parameters are optimized to reduce 2nd geometric factor.
 - No focusing in y direction in the PtP method.



Simulation parameters

Arc length of the dipole	0.5 m
Bending angle of the dipole	20 mrad
Strength of Q1 and Q4	0.33033 m^{-2}
Strength of Q2 and Q3	-0.91651 m^{-2}
Length of the quadrupoles	0.3 m
Frequency of the deflecting cavity	11.424 GHz
Length of the deflecting cavity	1.5 m
Voltage of the deflecting cavity	29 MV

Verification: ELEGANT simulation

Resolution sweep result

- Distance of chicane: 4 m
- $E = 4.3 \text{ GeV}$
- $\varepsilon_{project,n} = 0.6 \text{ mm} \cdot \text{mrad}$
- $T_{266} = 0.058 \text{ m}$, $T_{566} = 0.021 \text{ m}$
- $f = 12.4 \text{ m}$
- $M = 20$

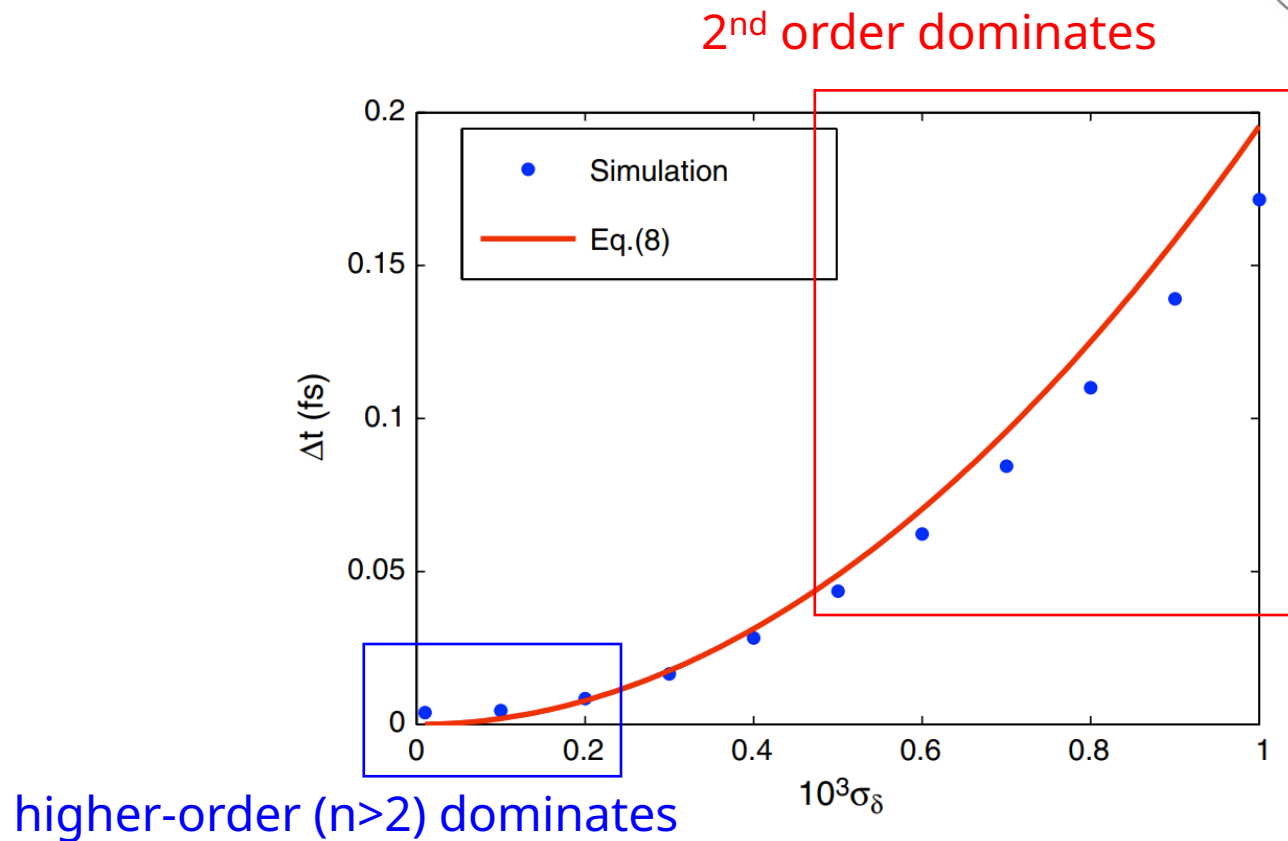


FIG. 2. Resolution in bunch length measurement: blue dots are the resolution determined from the beam size on the screen using ELEGANT simulation and the red curve is the results from Eq. (8).

$$\rightarrow \Delta t = 0.2 \text{ fs for } \sigma_p \sim 10^{-3}$$

Verification: Simulation using LCLS beam

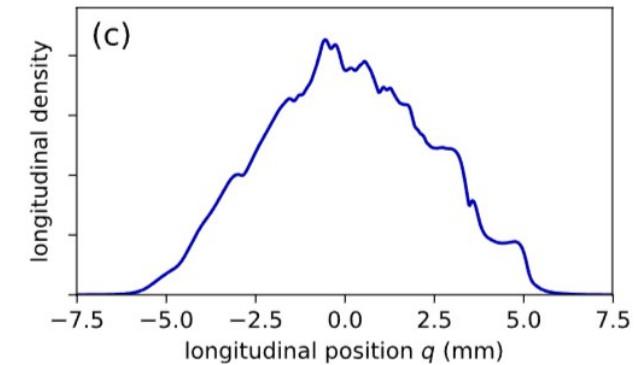
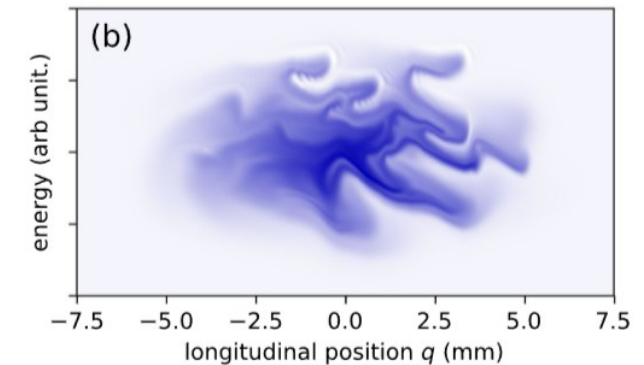
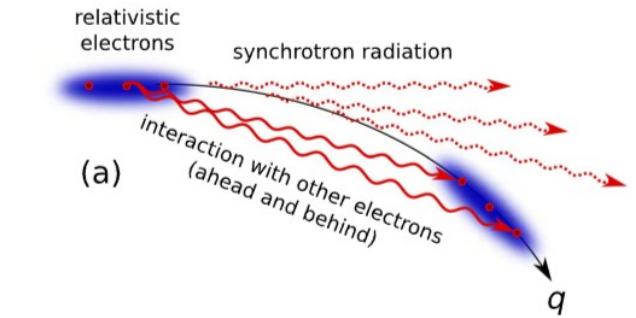
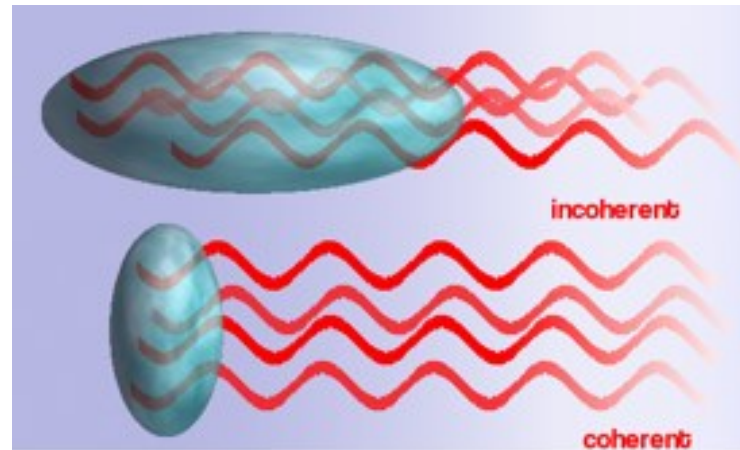
Coherent Synchrotron Radiation (CSR) effects

- SR emitted from the tail electrons catches up with the head electrons.
- Tail electrons lose energy due to CSR.
- Head electrons gain energy from the radiation.

→ Worse ε , causing micro-bunching instability (MBI)

$$\begin{bmatrix} x_c \\ x'_c \\ z_c \\ \delta_c \end{bmatrix} = \begin{bmatrix} -1 & R_{12} & 0 & \eta \\ 0 & -1 & 0 & R_{24} \\ 0 & -\eta & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} x_i \\ x'_i \\ z_i \\ \delta_i \end{bmatrix}$$

$$x_f = x_i + Lx'_i + kLz_i$$



Verification: Simulation using LCLS beam

Overcompressed beam: ISR and CSR effect

- Ground truth: 3.52 fs
- predicted: 3.07 fs w/ CSR, ISR

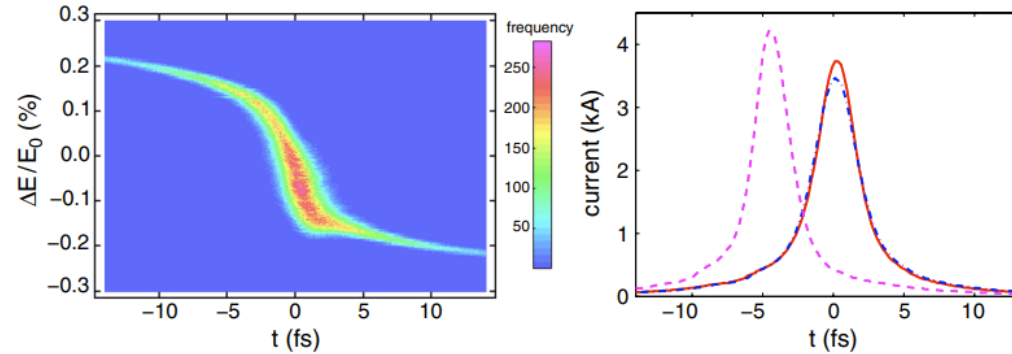


FIG. 3. (left) Longitudinal phase space for the overcompressed beam (bunch head to the left). (right) Temporal profiles: true temporal profile (red solid line); predicted temporal profile without considering ISR and CSR effect (blue dash-dotted line); predicted temporal profile with ISR and CSR taken into account (magenta dashed line).

$$\sigma_p = 1.1 \times 10^{-3}$$

$$\sigma_{t,res} \sim 0.2 \text{ fs}$$

Undercompressed beam: ISR and CSR effect

- Ground truth: 8.50 fs
- Predicted: 8.55 fs w/ CSR, ISR

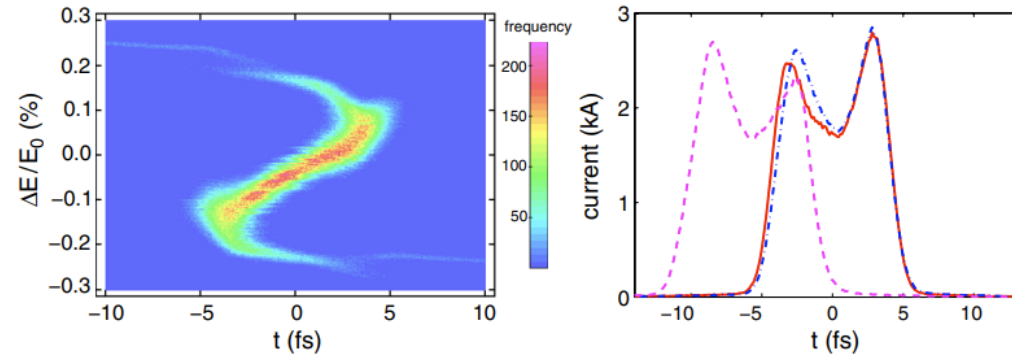


FIG. 4. (left) Longitudinal phase space for the undercompressed beam (bunch head to the left). (right) Temporal profiles: true temporal profile (red solid line); predicted temporal profile without considering ISR and CSR effect (blue dash-dotted line); predicted temporal profile with ISR and CSR taken into account (magenta dashed line).

Discussion

- Removing the core-beam chirp reduces global energy spread and improves resolution.
- With full chirp compensation (ideal LCLS case):
 - Energy spread ≈ 1.5 MeV.
 - Resolution improves to ~ 20 as.
- Increasing dipole spacing and reducing θ / k_{QM} helps mitigate 2nd order effects.

XTCAV optimization (2016)

		Non-SLED	SLED
Klystron power	MW	50	40
Klystron pulse length	μs	0.2	1
XTCAV kick amplitude	MV	45	85
Effective streak at 6 GeV	mm per °X-band	4.12	8.45
Temporal resolution IIXR (10kV)	fs	4	2
Temporal resolution SXR (1kV)	fs	1	0.5

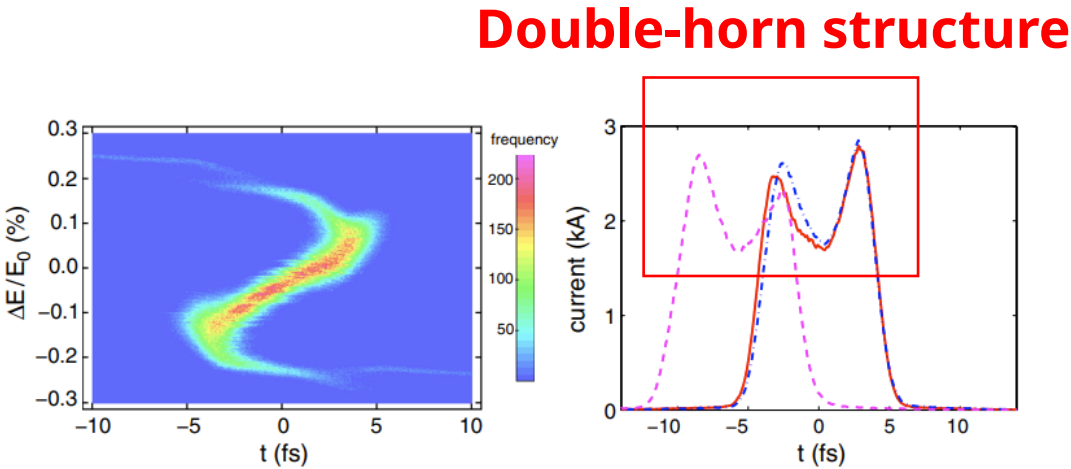


FIG. 4. (left) Longitudinal phase space for the undercompressed beam (bunch head to the left). (right) Temporal profiles: true temporal profile (red solid line); predicted temporal profile without considering ISR and CSR effect (blue dash-dotted line); predicted temporal profile with ISR and CSR taken into account (magenta dashed line).