가속기 자석 Part-II

전자석 평가를 위한 해석 기초

한가람 포항가속기연구소

2025-1 가속기실험실습 첨단원자력공학부/포항공과대학교

Field Quality



실제 자석은?



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Rotating Coil Magnetometers



8 mm

380 mm

Review of Generalized Magnetic Multipole Model

(2023/10/24-)

$$F = \vec{A} + iV$$
- 2D magnetic field model in complex plane의 당위성 증명
- 전류가 있을 때와 없을 때 Maxwell equatio을 만족하는지 평가 $i \frac{\partial F}{\partial x} = \frac{\partial F}{\partial y}$ - Cauchy Riemann 조건을 만족함을 보임
- 묵시적으로 conformal mapping이 가능함도 보임 $F = C_n z^n$ - 일반화된 멀티폴 포텐셜 모델링 $B_x = -\partial_x V$, $B_y = -\partial_x \vec{A}$ $B_x = \partial_y \vec{A}$, $B_y = -\partial_y V$ $B^* = i\partial_z F = iC_n(\partial_z z^n) = iC_n(nz^{n-1})$ - 일반화된 멀티폴 필드맵, 극좌표계에서 확인 $B^*_{n,main} \Big|_{max} = \frac{1}{\pi} \oint \left(\vec{B}(\theta) \Big|_{r=r_0} \cdot \hat{r} \right) \sin n\theta d\theta$ $C_{n,main} = -\frac{1}{nr_0^{n-1}} B^*_{n,main} \Big|_{max} [T/m^{n-1}]$ $B^*_{n,skew} \Big|_{max} = \frac{1}{\pi} \oint \left(\vec{B}(\theta) \Big|_{r=r_0} \cdot \hat{r} \right) \cos n\theta d\theta$ $C_{n,skew} = -\frac{1}{nr_0^{n-1}} B^*_{n,skew} \Big|_{max} [T/m^{n-1}]$



- Fundamental field component 에 정규화된 다극 성분 분석식

Modelling of 2D Magnetic Multipole : w/o current source 모델 당위성 증명

$$\mathbf{F} = \vec{A} + iV$$

If there's no current source (inside the closed loop)

$$\vec{\nabla} \times \vec{B} = 0 \qquad \qquad \vec{\nabla} \cdot \vec{B} = 0$$

Expressing \vec{B} by potential functions \vec{A} and V

$$\vec{B} = \vec{\nabla} \times \vec{A} = \begin{bmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{bmatrix}$$
$$\vec{B} = -\vec{\nabla}V = -\frac{\partial B}{\partial x}\hat{x} - \frac{\partial B}{\partial y}\hat{y} - \frac{\partial B}{\partial z}\hat{z}$$

Laplace equations, $\nabla^2 \vec{A} = 0$ and $\nabla^2 V = 0$

$$\vec{V} \times \vec{B} = \vec{\nabla} \times (\vec{\nabla} \times \vec{A})$$

= $\vec{\nabla} (\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A}$

$$\nabla^2 \vec{A} = 0 \quad (: \vec{\nabla} \cdot \vec{A} = 0) \quad (1)$$

(2)

$$\vec{\nabla} \cdot \vec{B} = -\vec{\nabla} \cdot \vec{\nabla} V$$
$$\nabla^2 V = 0$$

$$\therefore \nabla^2 \mathbf{F} = \nabla^2 \vec{A} + i \nabla^2 V = 0$$

 $F = \vec{A} + iV$ 로 모델링 하면, current source가 없을 때 "Gauss's low for magnetism" of Maxwell's Eq.를 만족함

Modelling of 2D Magnetic Multipole : with current source 모델 당위성 증명

Ampere's circuital law with DC current source can be written as

$$\vec{\nabla} \times \vec{B} = (\partial_y B_z - \partial_z B_y)\hat{x} + (\partial_z B_x - \partial_x B_z)\hat{y} + (\partial_x B_y - \partial_y B_x)\hat{z}$$
$$= \mu (J_x \hat{x} + J_y \hat{y} + J_z \hat{z})$$
(1)

If B is given $B = B_{\chi} + i B_{\gamma}$, then its conjugate form is



Equating (1) and (2), $\frac{\partial}{\partial x} \left(-\frac{\partial F}{\partial x} \right) - \frac{\partial}{\partial v} \left(\frac{\partial F}{\partial y} \right) = \mu J_z$

$$\nabla^2 F = \partial_x^2 F + \partial_y^2 F = -\mu J_z$$

 $F = \vec{A} + iV$ 로 모델링 하면, current source가 있을 때도 "Gauss's low for magnetism" of Maxwell's Eq.를 만족함 2025-02-28

Modelling of 2D Magnetic Multipole : $F = \vec{A} + iV$

odelling of 2D Magnetic Multipole :
$$F = \vec{A} + iV$$

 $E^{III} = i\partial_z F = i\left(\frac{\partial A}{\partial x} + i\frac{\partial V}{\partial x}\right) = -\frac{\partial V}{\partial x} - i\left(-\frac{\partial A}{\partial x}\right)$
 $\partial_z F = \frac{\partial}{\partial z}(A + iV) = \lim_{\Delta z \to 0} \frac{\Delta(A + iV)}{\Delta(x + iy)} = \lim_{\Delta z \to 0} \frac{\Delta(A + iV)/\Delta x}{1 + i\frac{\Delta y}{\Delta x}} = \frac{\partial A}{\partial x} + i\frac{\partial V}{\partial x}$
 $B_x = -\frac{\partial V}{\partial x}, \quad B_y = -\frac{\partial \vec{A}}{\partial x}$
 $B^* = i\partial_z F = i\left(-i\frac{\partial A}{\partial y} + \frac{\partial V}{\partial y}\right) = \frac{\partial A}{\partial y} - i\left(-\frac{\partial V}{\partial y}\right)$
 $\partial_z F = \frac{\partial}{\partial z}(A + iV) = \lim_{\Delta z \to 0} \frac{\Delta(A + iV)}{\Delta(x + iy)} = \lim_{\Delta z \to 0} \frac{\Delta(A + iV)/\Delta y}{\frac{\Delta x}{\Delta y} + i} = -i\frac{\partial A}{\partial y} + \frac{\partial V}{\partial y}$
 $B_x = -\frac{\partial \vec{A}}{\partial y}, \quad B_y = -\frac{\partial \vec{A}}{\partial y}$

 $F = \vec{A} + iV$ satisfies Cauchy Riemann condition $i \frac{\partial F}{\partial x} = \frac{\partial F}{\partial y}$, which enables conformal mapping

Multipoles : equipotential lines using F

일반화된 퍼텐셜 맵핑

2p, n=1
$$F = C_1 z$$

 $= C_1(x + iy) = A + iV$ $x = \frac{\vec{A}}{C_1}, \quad y = \frac{V}{C_1}$
4p, n=2 $F = C_2 z^2$ $x^2 - y^2 = \frac{\vec{A}}{C_2}, \quad 2xy = \frac{V}{C_2}$
6p, n=3 $F = C_3 z^3$ $x^3 - 3xy^2 = \frac{\vec{A}}{C_3}, \quad 3x^2y - y^3 = \frac{V}{C_3}$
8p, n=4 $F = C_4 z^4$ $x^4 - 6x^2y^2 + y^4 = \frac{\vec{A}}{C_4}, \quad 4x^3y - 4xy^3 = \frac{V}{C_4}$
10p, n=5 $F = C_5 z^5$ $x^5 - 10x^3y^2 + 5xy^4 = \frac{\vec{A}}{C_5}, \quad 5x^4y - 10x^2y^3 + y^5 = \frac{V}{C_4}$

2025-02-28

Multipoles : Field strength using F

일반화된 필드 맵핑

2p, n=1
$$x = \frac{\vec{A}}{C_1}, \quad y = \frac{V}{C_1},$$
 $B_x = -\partial_x V = 0, \quad B_y = -\partial_y V = -C_1$ 4p, n=2 $x^2 - y^2 = \frac{\vec{A}}{C_2}, \quad 2xy = \frac{V}{C_2},$ $B_x = -\partial_x V = -2C_2 y, \quad B_y = -\partial_x A = -2C_2 x$ 6p, n=3 $x^3 - 3xy^2 = \frac{\vec{A}}{C_3}, \quad 3x^2y - y^3 = \frac{V}{C_3}$ $B_x = \partial_y A = -6C_3 xy, \\ B_y = -\partial_y V = -3C_3 x^2 + 3C_3 y^2$ 8p, n=4 $x^4 - 6x^2y^2 + y^4 = \frac{\vec{A}}{C_4}, \quad 4x^3y - 4xy^3 = \frac{V}{C_4}$ \vdots 10p, n=5 $x^5 - 10x^3y^2 + 5xy^4 = \frac{\vec{A}}{C_5}, \quad 5x^4y - 10x^2y^3 + y^5 = \frac{V}{C_4}$ $\Box \Box \Delta A = -\partial_x V, \quad B_y = -\partial_x A$

2025-02-28

X-V-A, YA-V 암기용

 $B_y = -\partial_y V$

 $B_x = \partial_y \vec{A}$,

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Multipoles : field strength using F in polar coordinate

필드 맵핑 극좌표계

 $B^* = i\partial_z F = iC_n(\partial_z z^n) = iC_n(nz^{n-1})$

2p, n=1 $B_1^* = iC_1$ $B_{1y} = -C_1$

4p, n=2 $B_2^* = iC_2(2|z|e^{i\theta}) = i2C_2|z|(\cos\theta + i\sin\theta)$ $B_{2x} = -2C_2|z|\sin\theta$, $B_{2y} = -2C_2|z|\cos\theta$

6p, n=3 $B_3^* = iC_3(3|z|^2 e^{i2\theta}) = i3C_3|z|^2(\cos 2\theta + i\sin 2\theta)$ $B_{3x} = -3C_3|z|^2\sin 2\theta$, $B_{3y} = -3C_3|z|^2\cos 2\theta$

8p, n=4 $B_4^* = iC_4(4|z|^3 e^{i3\theta}) = i4C_4|z|^3(\cos 3\theta + i\sin 3\theta)$ $B_{4x} = -4C_4|z|^3\sin 3\theta$, $B_{4y} = -4C_4|z|^3\cos 3\theta$

Multipole Analysis

필드 맵핑 극좌표계

$$B^* = i\partial_z F = iC_n(\partial_z z^n) = iC_n(nz^{n-1})$$
$$B_{nx} = -nC_n|z|^{n-1}\sin(n-1)\theta$$
$$B_{ny} = -nC_n|z|^{n-1}\cos(n-1)\theta$$

Tr plane에서 중심 기준 r 고정, 한 바퀴 적분했을 때 orthogonality에 의해 아래와 같이 적분식으로 정리 가능

$$C_n = \frac{-1}{n\pi r_0^{n-1}} \oint B_x(\theta) \Big|_{r=r_0} \cdot \sin(n-1)\theta \, d\theta \qquad [T/m^{n-1}]$$

$$C_n = \frac{-1}{n\pi r_0^{n-1}} \oint B_y(\theta) \Big|_{r=r_0} \cdot \cos(n-1)\theta \, d\theta \qquad [T/m^{n-1}]$$

그런데 이렇게 쓰면 dipole skew만 볼 수 없음, 아래와 같이 바꿔 쓸 수 있음

$$C_{n,main} = \frac{-1}{n\pi r_0^{n-1}} \oint \left(\vec{B}(\theta) \Big|_{r=r_0} \cdot \hat{r} \right) \sin n\theta \ d\theta$$

$$C_{n,skew} = \frac{-1}{n\pi r_0^{n-1}} \oint \left(\vec{B}(\theta) \Big|_{r=r_0} \cdot \hat{r} \right) \cos n\theta \ d\theta$$

Multipole Analysis: multipole field error

(2023/10/25)

Normalized to the fundamental component
$$\frac{|B_n^*|}{|B_N^*|}$$
 is proportional to $\left|\frac{C_n}{C_N}\right|$
$$\frac{|B_n^*|}{|B_N^*|} = \frac{n}{N} r_0^{n-N} \left|\frac{C_n}{C_N}\right|$$
$$C'_n = C_n r_0^n \quad [T \cdot m] \qquad \left|\frac{C_n}{C_N}\right| = \left|\frac{C'_n}{C'_N}\right|$$

$$B_{n,main}^{*}\Big|_{max} = \frac{1}{\pi} \oint \left(\vec{B}(\theta)\Big|_{r=r_{0}} \cdot \hat{r}\right) \sin n\theta \, d\theta$$
$$B_{n,skew}^{*}\Big|_{max} = \frac{1}{\pi} \oint \left(\vec{B}(\theta)\Big|_{r=r_{0}} \cdot \hat{r}\right) \cos n\theta \, d\theta$$

$$C'_{n,main} = -\frac{r_0}{n} B^*_{n,main} \Big|_{max} [T \cdot m] \qquad C_{n,main} = -\frac{1}{nr_0^{n-1}} B^*_{n,main} \Big|_{max} [T/m^{n-1}] \\C'_{n,skew} = -\frac{r_0}{n} B^*_{n,skew} \Big|_{max} [T \cdot m] \qquad C_{n,skew} = -\frac{1}{nr_0^{n-1}} B^*_{n,skew} \Big|_{max} [T/m^{n-1}]$$

Code Implementation

```
double t Bmain, Bskew;
ThreeVector x probe;
double_t poles[13] = {1.0, 2.0, 3.0, 4.0, 5.0, 6.0, 7.0, 8.0, 9.0, 10.0, 11.0, 12.0, 13.0};
fo multipole << time*beta_e(p.GetKineticEnergy())*c;</pre>
for (const auto& mp : poles)
{
    Bmain = 0.; Bskew = 0.;
    fo_multipole << ",";</pre>
    for (const auto& th : ths)
    {
        ThreeVector dr = mp radius * x pr hat;
        dr.rotate(s hat, th);
        x probe = v0 + dr;
        ft->GetBField(x probe, e3b3);
        Bmain += ThreeVector(e3b3[3], e3b3[4], e3b3[5]).dot(dr.unit()) * sin(mp * th) * dth;
        Bskew += ThreeVector(e3b3[3], e3b3[4], e3b3[5]).dot(dr.unit()) * cos(mp * th) * dth;
        fo multipole scanpath.write((char*)&x_probe, sizeof(ThreeVector));
    Bmain /= pi; Bskew /= pi;
    double_t Cpr_main = -Bmain * mp_radius / mp, Cpr_skew = -Bskew * mp_radius / mp;
    double t Cmain = -Bmain / mp / pow(mp radius, mp-1), Cskew = -Bskew / mp / pow(mp radius, mp - 1);
    fo multipole
        << Bmain / tesla << "," << Bskew / tesla << ","
        << Cpr_main / (tesla * m) << "," << Cpr_main / (tesla * m) << ","
        << Cmain / (tesla / pow(m, mp - 1)) << "," << Cmain / (tesla / pow(m, mp - 1));
```

Artificial Tabulated Field-map Generation

1. Sector bend dipole

2. Sector bend dipole + quadrupole

1. Sector bend dipole

Specification

ROT_ANGLE = 6 * deg C1 = 1.4 * tesla BRHO = brho_e(4*GeV) RHO = BRHO/C1

Geometry computation

$$\begin{split} \vec{x} &= \{ \vec{x}_i \mid A \& B \& C \& D \& E \} \\ A: \quad D_{l1}(\vec{x}_i) > 0, \\ B: \quad D_{l2}(\vec{x}_i) < 0, \\ C: \quad R_{inner}(\vec{x}_i) > \rho - W_0, \\ D: \quad R_{outer}(\vec{x}_i) < \rho + W_0, \\ E: \quad |\vec{x}_i \cdot \hat{y}| < G_0 \end{split}$$

 $B(\vec{x}) = C_{1'}$

otherwise $B(\vec{x}) = 0$



1. Sector bend dipole (cont'd)



1. Sector bend dipole (cont'd), results



1. Sector bend dipole (cont'd), results



1. Sector bend dipole (cont'd), results



2. Sector bend dipole + quadrupole

$$\vec{B}'(\vec{x}') = (-2C_2y') \hat{x}' + (C_1 - 2C_2x') \hat{y}'$$

$$\vec{x}' = -(|\vec{r}| - \rho_0) \hat{x}' + y \hat{y}'$$

$$|\vec{r}| = |\vec{x} - \vec{0}_{curv}|$$

$$\hat{x}' = \mathbf{R}(\theta)\hat{x}, \quad \hat{z}' = \mathbf{R}(\theta)\hat{z}, \quad \hat{y}' = \hat{y}$$

$$\vec{B}'(\vec{x}') = (-2C_2y) \mathbf{R}(\theta) \hat{x} + (C_1 - 2C_2x') \hat{y}$$

$$= (-2C_2y)(\cos\theta \hat{x} - \sin\theta \hat{z}) + (C_1 + 2C_2(|\vec{x} - \vec{0}_{curv}| - \rho_0)) \hat{y}$$

$$\hat{x}' = \hat{x}'$$

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2(

2. Sector bend dipole + quadrupole (cont'd)

```
while x < xtab1:
  v_i = TwoVector(z, x)
  a = line_1.get_distance_to_point(v_i)
  b = line_2.get_distance_to_point(v_i)
  if a < 0 and b > 0 and not circ_ir.is_inside(v_i) and circ_or.is_inside(v_i):
     r = TwoVector(x,y) - v_curv_org
     theta = (-yhat2D).getTheta(r)
     shat_prime = zhat.rot_origin(yhat, theta)
     xhat_prime = xhat.rot_origin(yhat, theta) # equals to -rhat
     yhat_prime = yhat.new()
     x_prime = -(abs(r)-RHO)
     y_prime = y
                                               C1 = -1.4 * tesla
                                               C2 = -20 *
     Bx_prime = -2.0 * C2 * y_prime
                                               tesla/meter
     By_prime = C1 - 2.0 * C2 * x_prime
     B_prime = Bx_prime * xhat_prime + By_prime * yhat_prime
     Bx = B_{prime.x}
     By = B_{prime.y}
     Bz = B_{prime.z}
  else:
     Bx, By, Bz = 0, 0, 0
```





2. Sector bend dipole + quadrupole (cont'd), results, r=20 mm



2. Sector bend dipole + quadrupole (cont'd), results , r=20 mm



2. Sector bend dipole + quadrupole (cont'd), results , r=2 mm, 20 mm

$$\frac{B_{n,main}^{*}}{\max} = \frac{1}{\pi} \oint \left(\vec{B}(\theta) \Big|_{r=r_{0}} \cdot \hat{r} \right) \sin n\theta \, d\theta$$

